

Lecture 18

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1 Overview

In this first lecture, we will learn how to compute simple Markov chains using resistor network.

2 Simple Markov Chains

A simple Markov chain is time-reversible, and thus can be represented by undirected graph. A formal definition of time reversible is stated as follows:

Definition 1. A Markov chain is time-reversible if there exists a stationary distribution π s.t.

$$\forall i, j \in S, \pi_i P_{ij} = \pi_j P_{ji}$$

Here, π_i, π_j is the probability at state i and j , respectively. P_{ij} is the transition probability from $i \rightarrow j$ when being at state i , and vice versa. A schematic illustration of a two-state simple Markov chain is shown below:

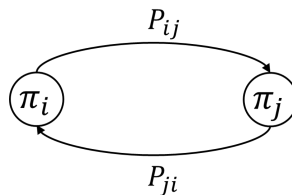


Figure 1: state diagram of simple two-state Markov chain

3 Representing Simple Markov Chain as Undirected Graph

For a simple Markov chain as defined above, we can use an undirected graph to represent it (i.e. we can use only one edge to connect two nodes). The **edge weight** for $e(i, j)$ is defined by:

$$C_{ij} = \pi_i P_{ij} = \pi_j P_{ji}$$

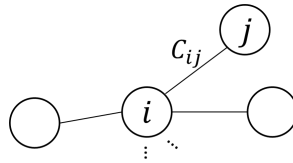


Figure 2: undirected graph representing node i and its surrounding nodes

Notice that $\sum_{j \in S(i)} C_{ij} = \sum_{j \in S(i)} \pi_i P_{ij} = \pi_i \sum_{j \in S(i)} P_{ij} = \pi_i \times 1 = \pi_i$, we can write P_{ij} as:

$$P_{ij} = \frac{C_{ij}}{\pi_i} = \frac{C_{ij}}{\sum_{j \in S(i)} C_{ij}}$$

Here, $S(i)$ denotes the space of all surrounding nodes of node i .

Note:

1. C_{ij} can be multiplied by a constant.
2. In general, $P_{ij} \neq P_{ji}$.
3. π_i is proportional to $\sum_{j \in S} C_{ij} = \text{deg}(i)$
4. In this lecture, we choose $C_{ij} = 1$ for all $(i, j) \in E$. As a result, $\pi_i = \frac{\text{deg}(i)}{2|E|}$, where $|E|$ denotes the total number of edges in this graph.

4 Compute Probability/Expectation of Markov Chains

4.1 Problem Description

Definition 2. Define the following symbols as:

$$\gamma_i = \Pr[\text{event happens, if we start random walk at state } i]$$

Note: Here "event" refers to "reach a certain state" or "go through an edge" or other similar events.

In the following texts, we will show how to computer γ_i .

4.2 Solve the Problem Using Markov Property

The idea is to use Markov property and linearity of expectations.

Assuming the event cannot take place within one step. For node i , we have:

$$\begin{aligned}
\gamma_i &= \sum_{j \in \mathcal{S}(i)} Pr[event|j] \times Pr[state(j)|state(i)] \\
&= \sum_{j \in \mathcal{S}(i)} \gamma_j P_{ij} \\
&= \frac{\sum_{j \in \mathcal{S}(i)} \gamma_j}{deg(i)}
\end{aligned}$$

If the graph has n nodes, we can write out n linear equations, i.e.

$$\gamma = P\gamma, \quad \gamma = (\gamma_1, \gamma_2, \dots, \gamma_n)^T$$

There are two kinds of uncertainties for the solution of the above linear system:
 If γ is a solution, then (1) $C \cdot \gamma$ is also a solution. (2) $\gamma + \lambda \vec{1}$ is also a solution.

4.3 Solve the Problem Using Resistor Network

Recall that in a resistor network, the voltage at node i is given by:

$$v_i = \arg \min_v \sum_{j \in \mathcal{S}(i)} \frac{1}{2} (v_j - v)^2$$

Here, $\frac{1}{2}(v_j - v)^2$ is proportional to the energy generated by resistors between node i and j .
 We can solve v_i by taking the derivative of this quadratic equation, and the result will be:

$$v_i = \frac{\sum_{j \in \mathcal{S}(i)} v_j}{deg(i)}$$

We can see that the expression of v_i is the same as the expression of γ_i in the previous section. Therefore, by solving the voltages in a resistor network, we can solve the event probability in a simple Markov chain.

5 Examples

5.1 Example 1

Among the n states, pick two states a and b .

The event is defined as "reaching a before reaching b"

And let $\gamma_i = Pr$ [event happens, if we start random walk at state i]

Solution:

In addition to the linear system equations $\gamma = P\gamma$, we notice that there are two boundary conditions:

$$\gamma_a = 1, \quad \gamma_b = 0$$

The boundary condition is analogous to applying a voltage source from a to b . In this way, we resolved the two uncertainties as described before, and we can get a unique solution.

5.2 Example 2

Among the n states, pick two states a and b .

The event is defined as "reaching a before returning to b "

And let $\gamma_i = Pr[\text{event happens, if we start random walk at state } i]$

The **escape probability** is defined as γ_b , i.e. the probability of reaching a before returning to b , if we start random walk at state b . Then the **escape probability** is given by:

Solution:

Notice that compared with Example 1, we only have one boundary condition $\gamma_a = 1$, which is not enough to solve the linear system.

But we can use the resistor network to solve γ_b .

$$\begin{aligned} \Pr[\text{escaping}] = \gamma_b &= \frac{\sum_{j \in S(b)} \gamma_j}{\text{deg}(b)} \\ &= \frac{1}{\text{deg}(b) \cdot R_{eff}} \end{aligned}$$

Here, $R_{eff} = \frac{1}{\text{current that goes into node } b}$ is the effective resistance between node a and node b .

5.3 Example 3

Among the n states, pick two states a and b .

Define the event as "get to state a ".

Define X_i as: $X_i = E[\text{number of steps to event (get to state } a), \text{ if starting from } i]$

Solution:

Using the same idea as in Section 4.2, we can get:

$$X_i = \frac{\sum_{j \in S(i)} X_j}{\text{deg}(i)} + 1$$

And we have one extra equation: $X_a = 0$.

As a result, we can solve the linear equations with unique solution. (Note: In comparison to Example 1 and 2, there is only one uncertain here. This is because in the equation above, there is an extra +1)